When the Born rule is not enough

Part I: Foreseeing future outcomes of measurements on one system

8

Part II: Repeated measurements on Unruh-DeWitt detectors in (3+1) Minkowski spacetime Quantum Hiking 2022

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Guillermo Garc**í**a-Pérez,

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Summary

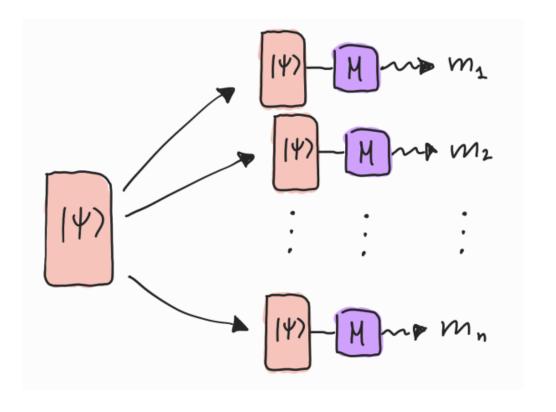
Part I

- Born rule
- Repeated Measurements (RM)
- Almost-Born probabilities

Part II

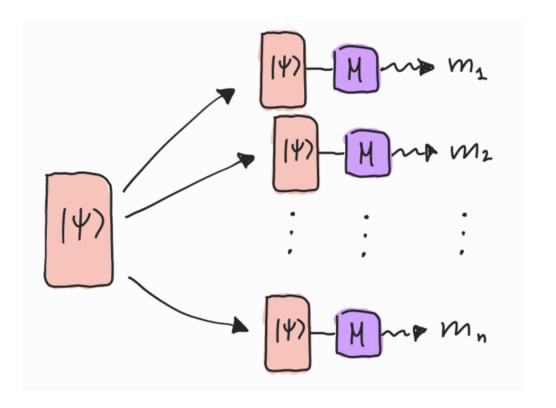
- Unruh-DeWitt detectors (UDW)
- RM on UDW
- Almost-Born rule for UDW

$$(\hat{E}_i, m_i)$$
 $|\psi\rangle$



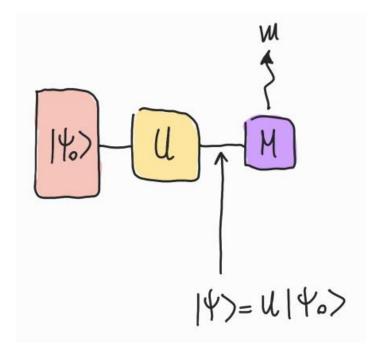
$$|\psi\rangle$$

$$\frac{N(m_k)}{\sum_i N(m_i)} \longrightarrow p_{m_k} = \langle \psi | \hat{E}_k | \psi \rangle$$



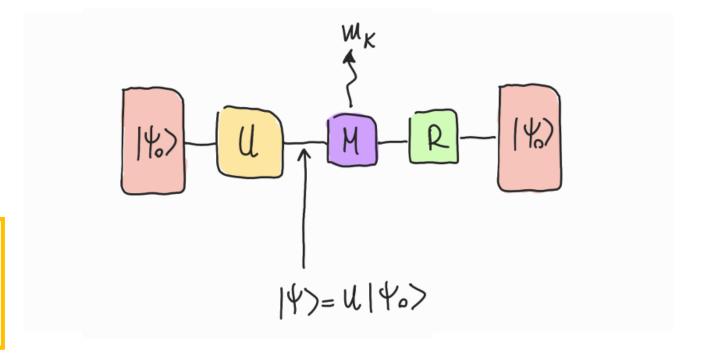
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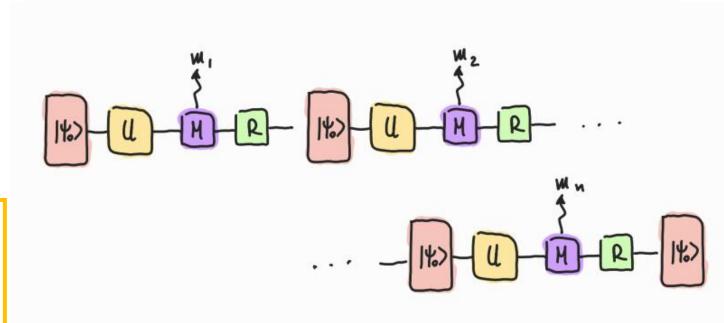
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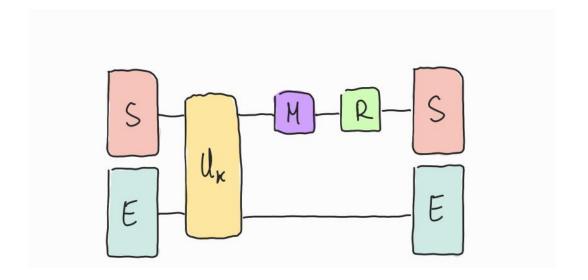


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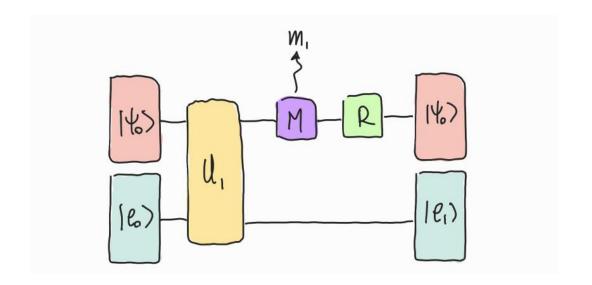


$$\{\mathcal{U}_1, \dots, \mathcal{U}_n\}$$
$$|g\rangle \otimes |f\rangle \to |g\rangle \otimes \left(\hat{V}_m(k)[f]|f\rangle\right)$$
$$\hat{V}_{m_k}(k)[f] = \frac{\langle m_k|\hat{\mathcal{U}}_k|g\rangle}{\sqrt{p_{m_k}(k)[f]}}$$



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$$|\phi\rangle = |\psi\rangle \otimes |e_0\rangle$$

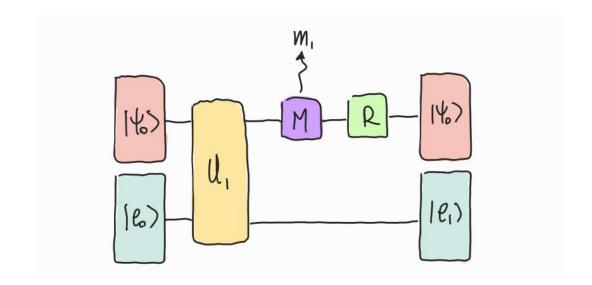


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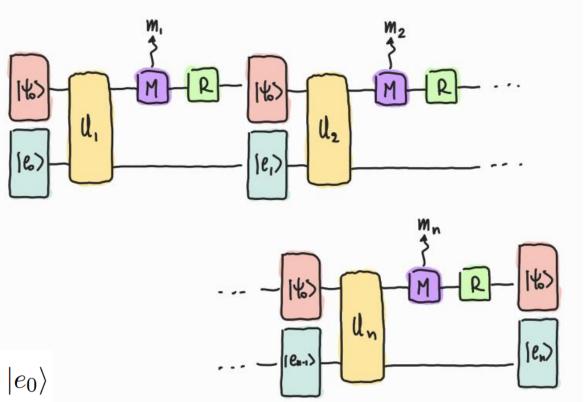


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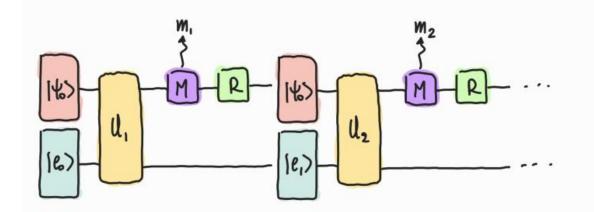


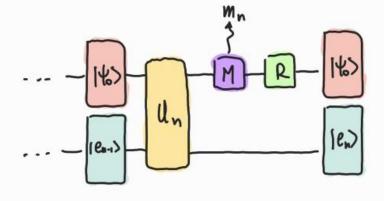
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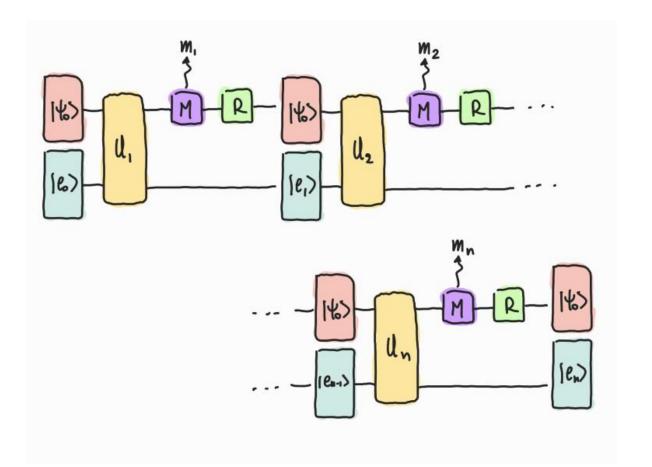




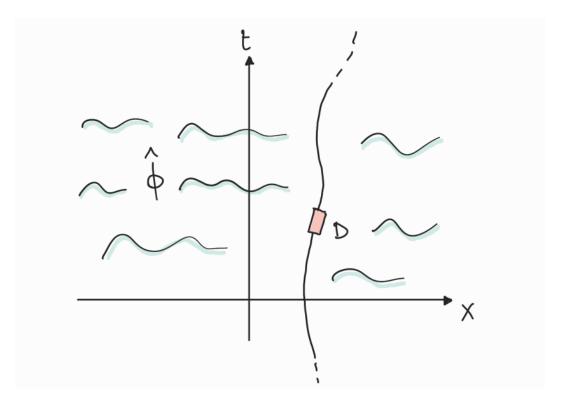
$$\frac{N(m_k)}{\sum_i N(m_i)} \implies p_{m_k} = \langle \psi | \hat{E}_k | \psi \rangle$$

Almost-Born probabilities

$$\hat{\mathcal{U}}_k = \hat{U} \otimes \mathbb{I} + \epsilon \sum_l \hat{A}_l \otimes \hat{B}_l(k) + \epsilon^2 \sum_l \hat{C}_l \otimes \hat{D}_l(k) + O(\epsilon^3)$$

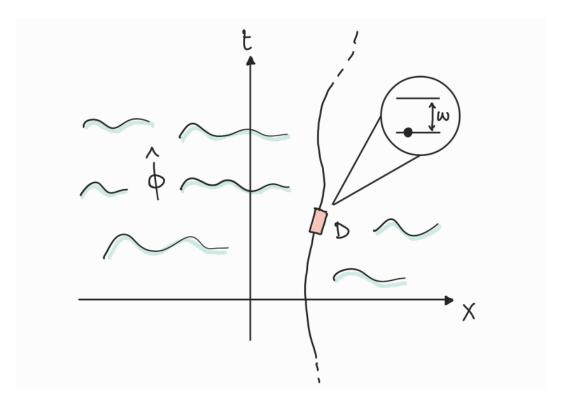


$$p_m(k)[f] = p_m + \epsilon Q_m^{(1)}(k)[f] + \epsilon^2 Q_m^{(2)}(k)[f] + O(\epsilon^3)$$



$$S = D \cup \phi \Rightarrow \mathcal{H} = \mathcal{H}_D \otimes \mathcal{H}_{\phi}$$

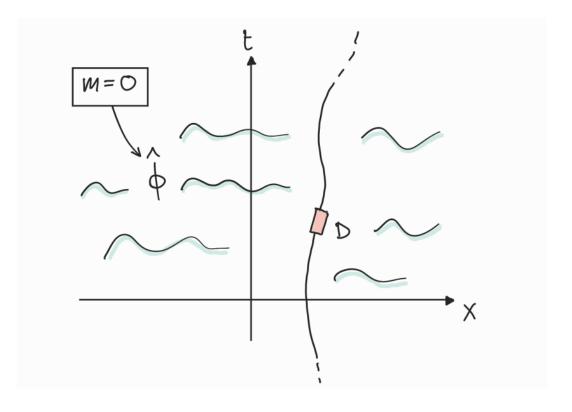
$$X(\tau) = (t(\tau), \mathbf{x}(\tau)), \ \hat{H}_D = \omega |e\rangle \langle e|$$



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$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi \implies \hat{H}_{\phi} = \int \left(\frac{1}{2} \hat{\pi}^2 + \frac{1}{2} (\nabla \hat{\phi})^2 \right)$$

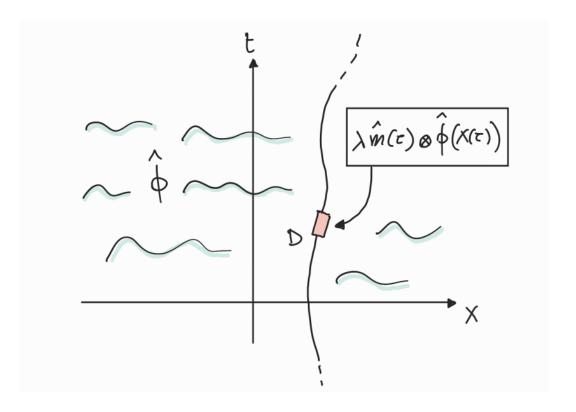


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$$\hat{H}_{int}(\tau) = \lambda \hat{m}(\tau) \otimes \hat{\phi}(X(\tau))$$

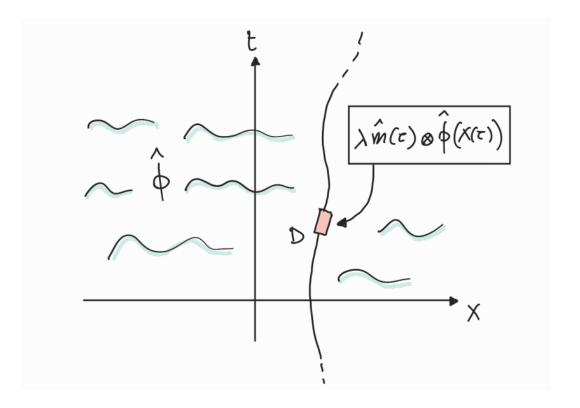


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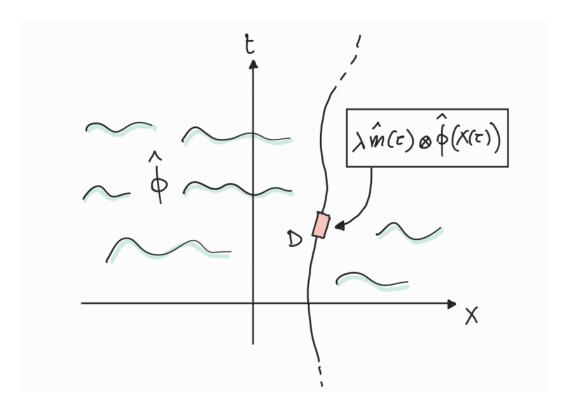
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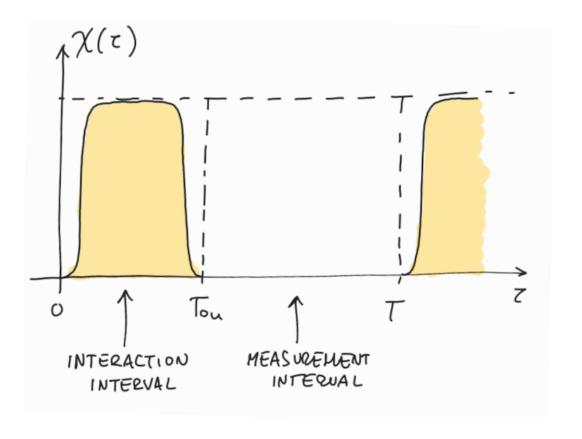
$$\hat{H}_{tot}(\tau) = \hat{H}_D \otimes \mathbb{I}_{\phi} + \mathbb{I}_{\phi} \otimes \hat{H}_{\phi} + \hat{H}_{int}(\tau)$$



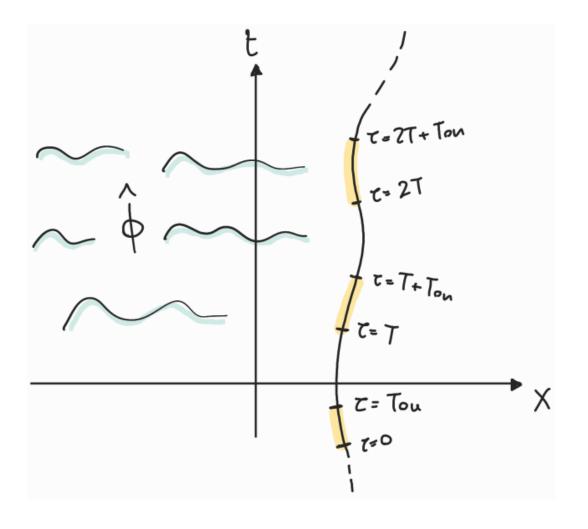
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The switching function

$$\chi(\tau) = \sum_{j} \chi_{j}(\tau) , \quad supp\{\chi_{j}(\tau)\} \subseteq [jT, jT + T_{on}]$$

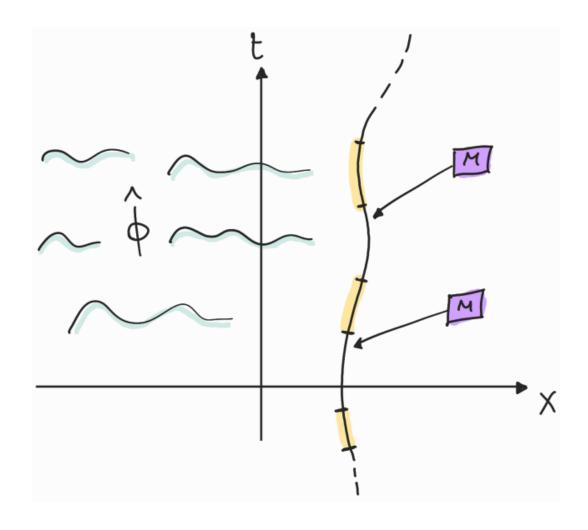


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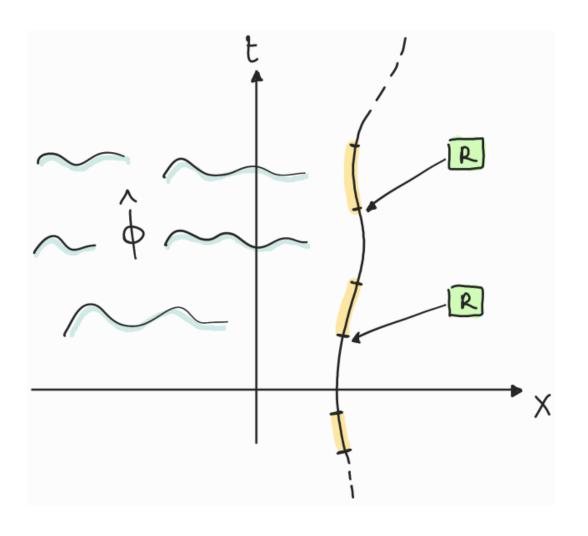
$$\begin{cases} \hat{M}_{0} = |g\rangle \langle g| \\ \hat{M}_{1} = |e\rangle \langle e| \end{cases}$$



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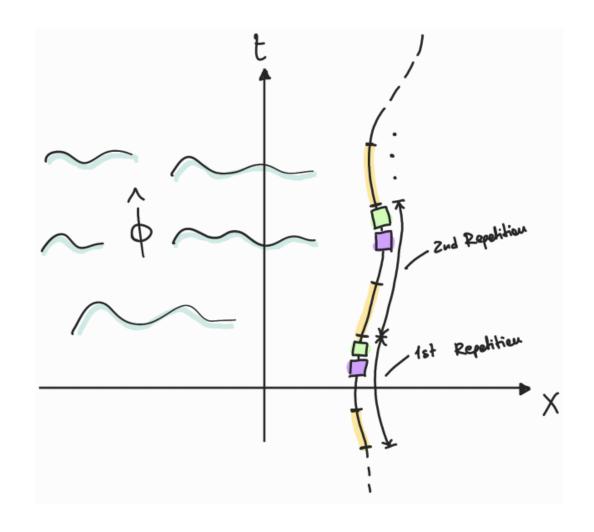
$$|\psi\rangle \xrightarrow{R} |g\rangle$$

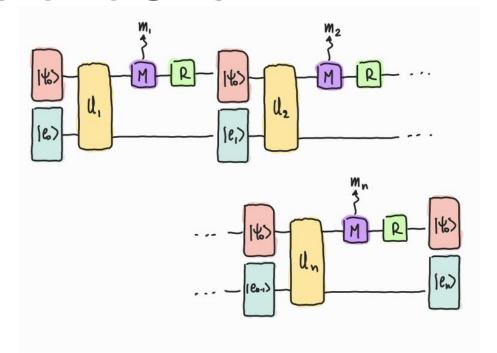


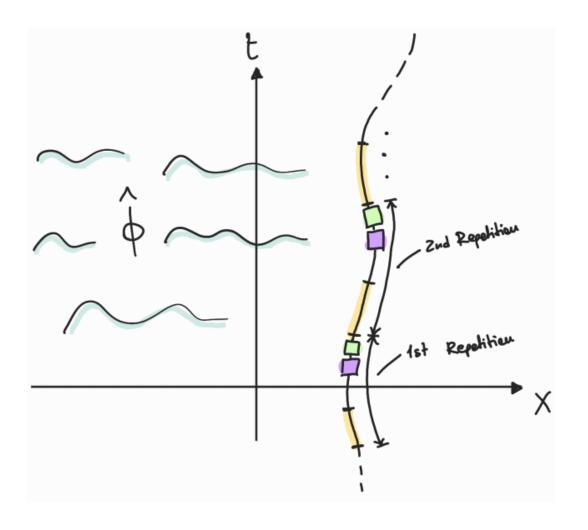
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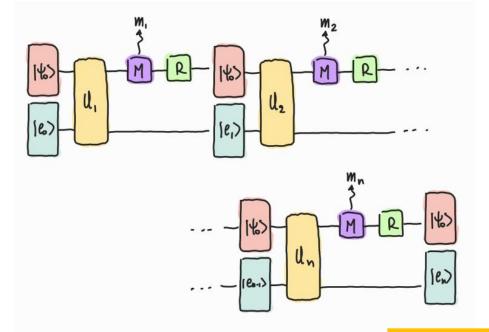
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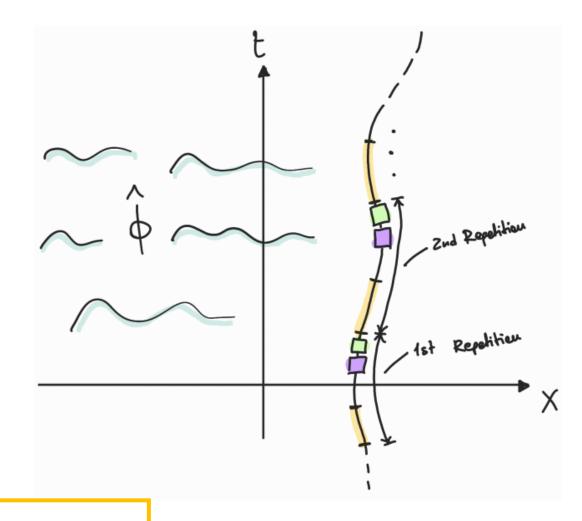
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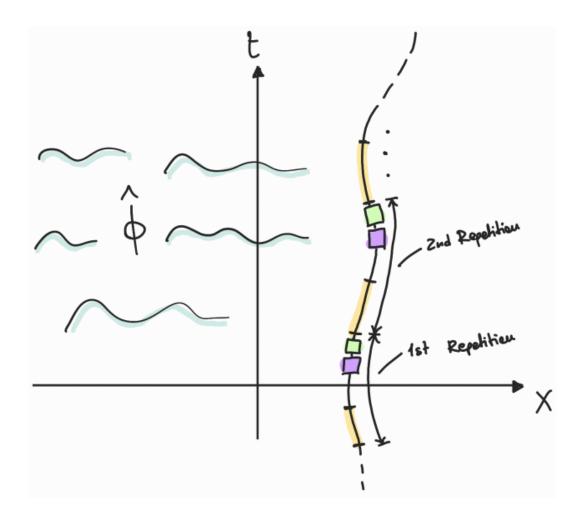


$$M_L=(m_1,\ldots,m_L)$$

$$M_L = (m_1, \dots, m_L)$$

$$M_L \iff (L; i_1, \dots, i_n)$$

$$(L; -) \equiv (0, \dots, 0)$$

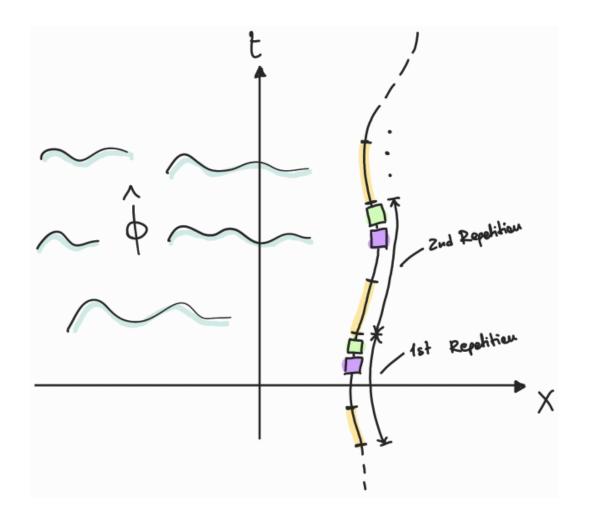


$$M_L = (m_1, \dots, m_L)$$

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$$M_7 = (0, 1, 0, 0, 0, 0, 1) \iff (7; 2, 7)$$



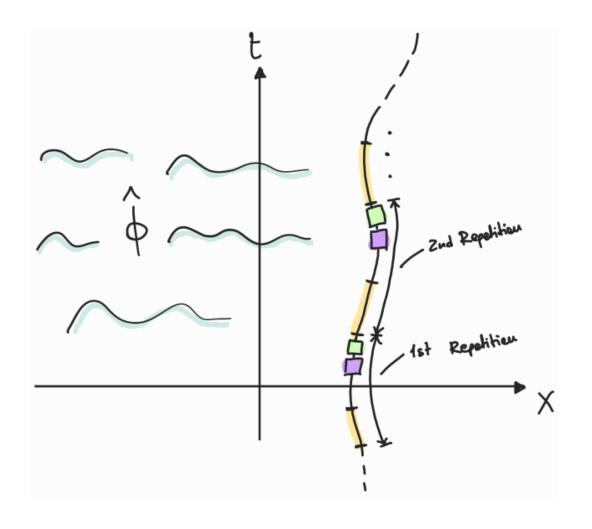
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$$P(M_L) = P(m_L|M_{L-1})P(M_{L-1})$$



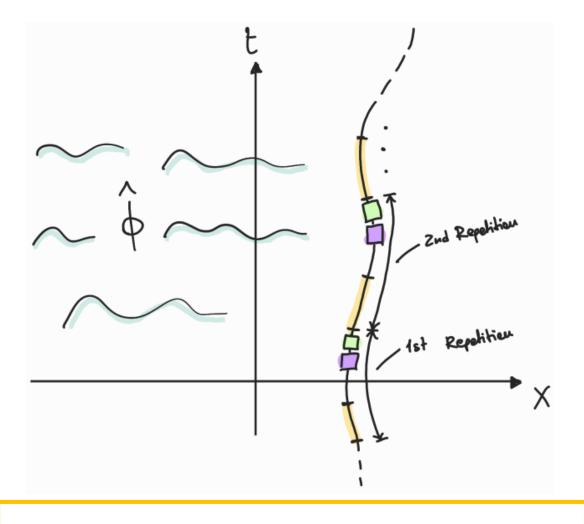
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$$P[0|(L; i_1, \dots, i_n)]$$
 $P[1|(L; i_1, \dots, i_n)]$

$$|g\rangle\otimes|0_{M}\rangle$$

$$|g\rangle\otimes|0_{M}\rangle-i\lambda\int_{LT}^{LT+T_{on}}d\tau\chi(\tau)\hat{m}(\tau)|g\rangle\otimes\hat{\phi}(X(\tau))|0_{M}\rangle+O(\lambda^{2})\sim|g\rangle\otimes|0_{M}\rangle+|e\rangle\otimes|\phi_{1}\rangle_{L}$$

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$$P[1|(L-1;-)] = \lambda^2 \int_{LT}^{LT+T_{on}} d\tau \int_{LT}^{LT+T_{on}} d\tau' \chi(\tau) \chi(\tau') e^{-i\omega(\tau-\tau')} \mathcal{W}_2(X(\tau'), X(\tau))$$

$$\mathcal{W}_2(X(\tau'), X(\tau)) = \langle 0_M | \hat{\phi}(X(\tau')) \hat{\phi}(X(\tau)) | 0_M \rangle$$

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$$|g\rangle\otimes|0_M\rangle$$
 \xrightarrow{R} $|g\rangle\otimes|\phi_1\rangle_L$

$$|g\rangle \otimes |\phi_1\rangle_{i_1} - i\lambda \int_{LT}^{LT+T_{on}} d\tau \chi(\tau) \hat{m}(\tau) |g\rangle \otimes \hat{\phi}(X(\tau)) |\phi_1\rangle_{i_1} + O(\lambda^2) \sim |g\rangle \otimes |0_M\rangle + |e\rangle \otimes |\phi_2\rangle_{L}$$

$$|g\rangle\otimes|\phi_1\rangle_{i_1}-i\lambda\int_{LT}^{LT+T_{on}}d\tau\chi(\tau)\hat{m}(\tau)|g\rangle\otimes\hat{\phi}(X(\tau))|\phi_1\rangle_{i_1}+O(\lambda^2)\sim|g\rangle\otimes|0_M\rangle+|e\rangle\otimes|\phi_2\rangle_{L}$$

$$P[1|(L-1;i_1)] = \frac{\lambda^4}{q} \iint_L \langle \phi_1 |_L \hat{\phi}(X(\tau_L')) \hat{\phi}(X(\tau_L)) | \phi_1 \rangle_L$$

= $\frac{\lambda^4}{q} \iint_L \iint_{i_1} W_4(X(\tau_2'), X(\tau_1'), X(\tau_1)X(\tau_2))$

•
$$\mathcal{W}_4(X(\tau_2'), X(\tau_1'), X(\tau_1)X(\tau_2)) = \langle 0_M | \hat{\phi}(X(\tau_{i_1}')) \hat{\phi}(X(\tau_L')) \hat{\phi}(X(\tau_L)) \hat{\phi}(X(\tau_{i_1})) | 0_M \rangle$$

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$$|\phi_1\rangle_{i_1} \to |\phi_2\rangle_L$$

$$P[1|(L-1;i_{1},...,i_{n-1})] =$$

$$= \frac{\lambda^{2n}}{q \times q(2;i_{2})...q(n-1;i_{n-1})} \iint_{L} \prod_{j=1}^{n-1} \iint_{j} \langle \phi_{n-1}|_{L} \hat{\phi}(X(\tau'_{L})) \hat{\phi}(X(\tau_{L})) | \phi_{i_{n-1}} \rangle_{L}$$

$$= \frac{\lambda^{2n}}{q \times q(2;i_{2})...q(n-1;i_{n-1})} \iint_{L} \prod_{j=1}^{n-1} \iint_{j} W_{2n}(X(\tau'_{L}), X(\tau_{L}),...)$$

Upper and Lower Bounds

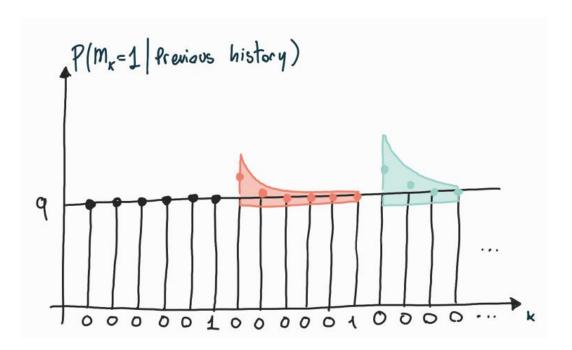
Using the **strong Huygens principle**, and under the assumption that

$$T_{on}\omega \leq \frac{\pi}{2}$$

$$T_{on} \ll T_{off}$$

One finds

$$\begin{cases} q \leq P[1|(L; i_1)] \leq q(1 + 2\gamma^2) \\ \frac{q}{(1 + 2\gamma^2)} \leq P[1|(L; i_1, i_2)] \leq q(1 + 6\gamma^2 + 8\gamma^3) \\ \dots \\ \frac{q}{(1 + \dots + \frac{(2n-3)!!}{\sqrt{e}}\gamma^{n-1})} \leq P[1|(L; i_1, \dots, i_{n-1})] \leq q(1 + \dots + \frac{(2n-1)!!}{\sqrt{e}}\gamma^n) \end{cases}$$



Upper and Lower Bounds

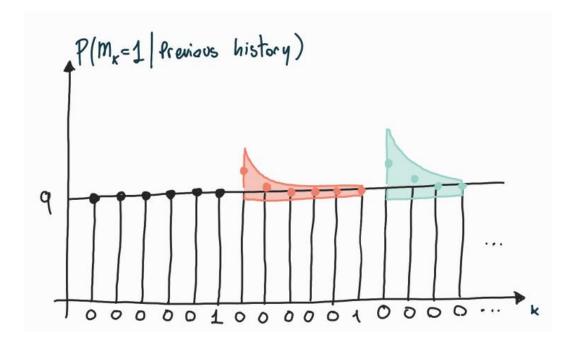
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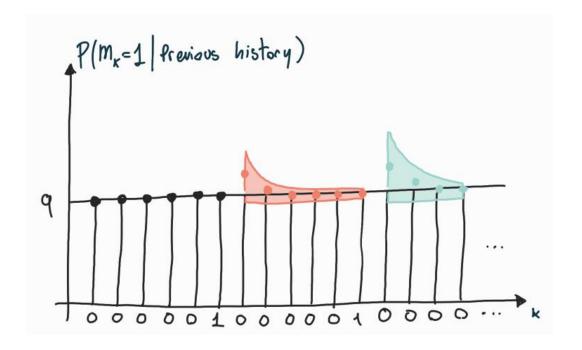
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$$p[1|(L; i_1, \dots, i_n)] \simeq q$$
$$p(M_L) \simeq q \times q^n (1-q)^{L-n}$$

Conclusions

- In some cases, even if holding the Born rule has no predicting power
- In weak RM one recovers an almost-Born rule
- FAPP, UDW detectors see strings consistent with asymptotic populations (for any trajectory) even in the RM scenario
- ...and more!